## 10.1. Spin- $\frac{1}{2}$ triangle

Consider three spins forming a triangle. They are described by the Hamiltonian

$$\mathcal{H} = J(\hat{\mathbf{S}}_1\hat{\mathbf{S}}_2 + \hat{\mathbf{S}}_1\hat{\mathbf{S}}_3 + \hat{\mathbf{S}}_2\hat{\mathbf{S}}_3)$$

where  $S_1 = S_2 = S_3 = \frac{1}{2}$  and J > 0. Solve this Hamiltonian similar to the spin- $\frac{1}{2}$  dimer shown in lecture 21.

(a) Construct the basis set. How many states does it contain, and how many eigenstates of  $\mathcal{H}$  do you expect?

(b) Show that  $\mathcal{H}$  commutes with  $\hat{S}^z = \hat{S}_1^z + \hat{S}_2^z + \hat{S}_3^z$ . You can then use  $S^z$  as a quantum number to distribute your basis states into several groups.

(c) Write down the block-diagonal matrix of  $\mathcal{H}$ .

(d) Diagonalize the matrix and obtain the eigenstates. Label them according to their  $S^z$  and  $S^2$  values where  $\hat{\mathbf{S}}^2 = (\hat{\mathbf{S}}_1 + \hat{\mathbf{S}}_2 + \hat{\mathbf{S}}_3)^2$ .

(e) Compare the result to the spin- $\frac{1}{2}$  dimer shown in the lecture. Why is the ground state non-degenerate in the case of the dimer and degenerate in the case of the triangle?

## 10.2. Classical version of spin waves

Obtain spin-wave excitations without quantum mechanics. Consider a chain of spins S with the exchange coupling J.

(a) Use the energy from the classical version of the Heisenberg model,

$$E = J \sum_{p} \mathbf{S}_{p} \mathbf{S}_{p+1}$$

to obtain the energy associated with spin p,

$$E_p = -g\mu_B \mathbf{S}_p \mathbf{B}_p^{\text{eff}}$$

where g = 2 and  $\mathbf{B}_p^{\text{eff}}$  is the effective magnetic field acting on  $\mathbf{S}_p$ .

(b) Use the equation of motion,

$$h\frac{d\mathbf{S}_p}{dt} = g\mu_B \mathbf{S}_p \times \mathbf{B}_p^{\text{eff}},$$

to derive the spin-wave excitation. To this end, assume that  $S_p^x$  and  $S_p^y$  are small, whereas  $S_p^z \simeq S$  (the excitation is a small tilt of the spins away from the z-axis). Write the equations of motion for  $S_p^x$  and  $S_p^y$  assuming the ferromagnetic ground state of the spin chain (J < 0).

(c) Use the ansatz

$$S_p^x = u e^{i(pqa-\omega t)}, \qquad S_p^y = v e^{i(pqa-\omega t)}$$

to solve the equations from (b). Here, u and v are the amplitudes, q is the wave number (1D propagation vector), and a is the lattice parameter. Show that u = -iv (hence the  $\pi/2$  shift between the oscillations of the x- and y-spin components) and obtain the dispersion relation for the ferromagnetic chain.

(d) Repeat the steps (b) and (c) for an antiferromagnet (J > 0). In this case,  $S_p^z = (-1)^p S$ , and you should write four equations of motion, two for  $S_p^x$  and  $S_p^y$  and two more for  $S_{p+1}^x$  and  $S_{p+1}^y$ , respectively. Combine them into two equations for  $S_p^+$  and  $S_{p+1}^+$  (the  $\pi/2$  shift between  $S_p^x$  and  $S_p^y$  is now implicitly taken into account within  $S_p^+ = S_p^x + iS_p^y$ ).

## Problem sheet 10: Collective spin states

(e) Now solve these equations using the ansatz

 $S_p^+ = u e^{i(pqa-\omega t)}, \qquad S_{p+1}^y = v e^{i((p+1)qa-\omega t)}$ 

Obtain the dispersion relation for an antiferromagnet and compare it to the ferromagnetic case.